

# The Hellan–Herrmann–Johnson and TDNNS method for nonlinear Koiter and Naghdi shells

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Michael Neunteufel (TU Wien)

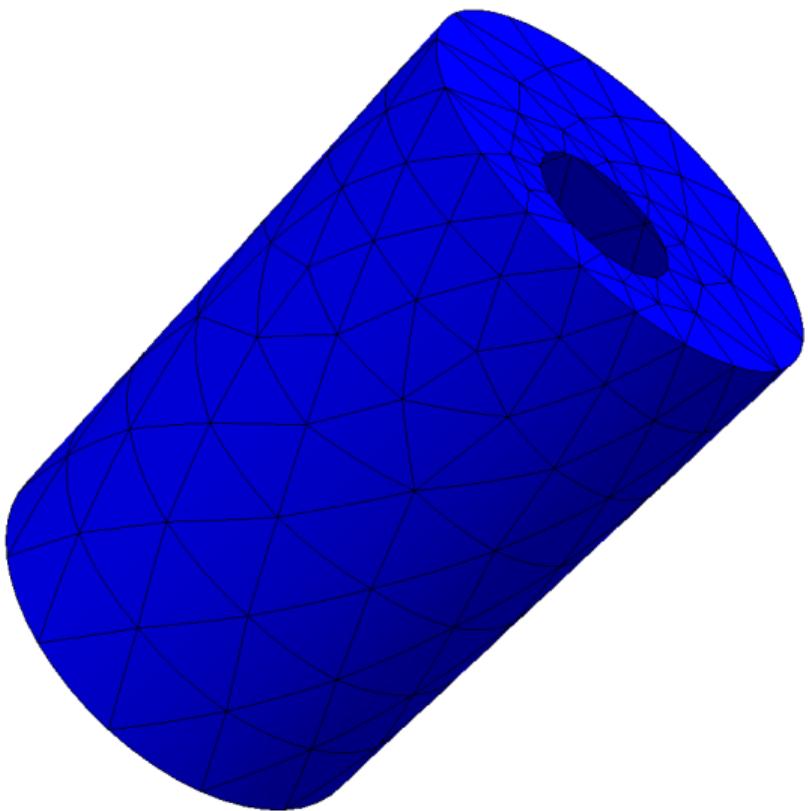
Joachim Schöberl (TU Wien)



**FWF** Austrian  
Science Fund



10th GACM Colloquium on Computational Mechanics, Vienna, Sep 11th, 2023









Nonlinear shells

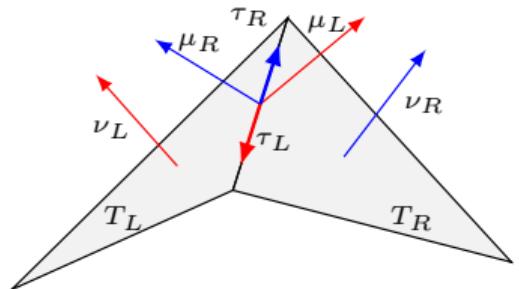
Membrane locking

Numerical examples

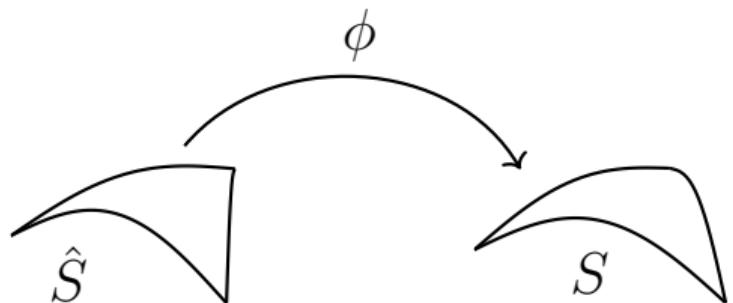
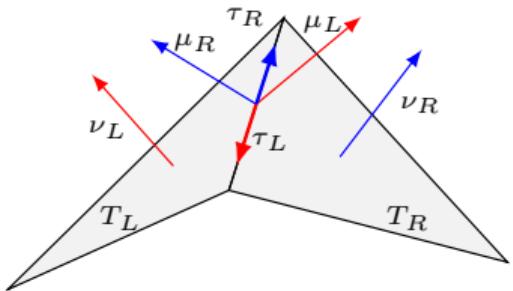
## Nonlinear shells

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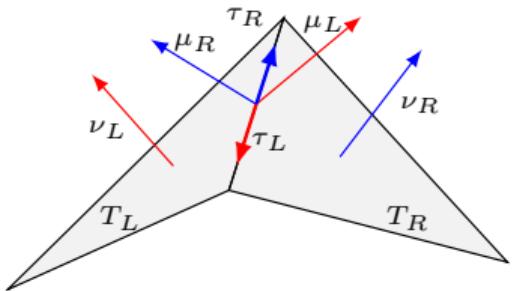
- Normal vector  $\nu$
- Tangent vector  $\tau$
- Element normal vector  $\mu = \nu \times \tau$



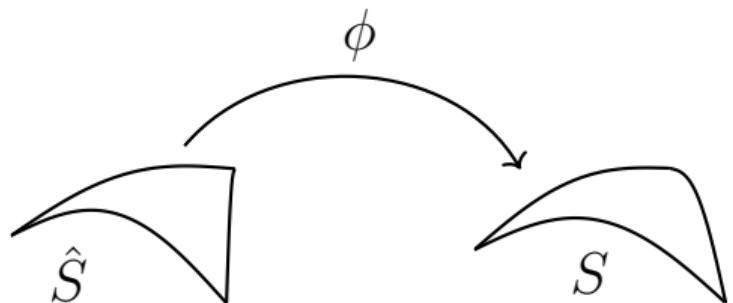
- Normal vector  $\hat{\nu}$
- Tangent vector  $\hat{\tau}$
- Element normal vector  $\hat{\mu} = \hat{\nu} \times \hat{\tau}$



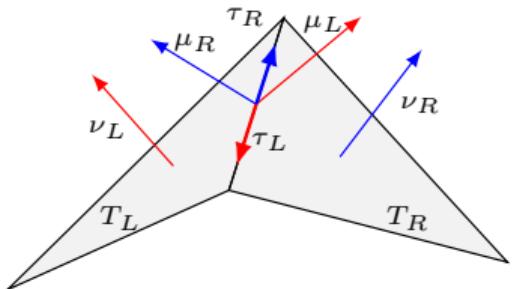
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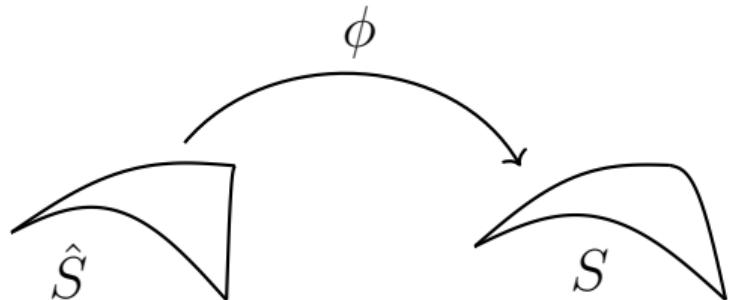
- $\mathbf{F} = \nabla_{\hat{\tau}} \phi$ ,  $J = \sqrt{\det(\mathbf{F}^\top \mathbf{F})}$



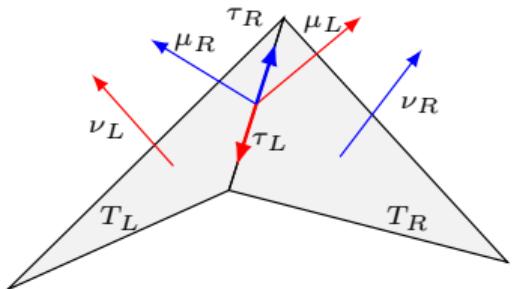
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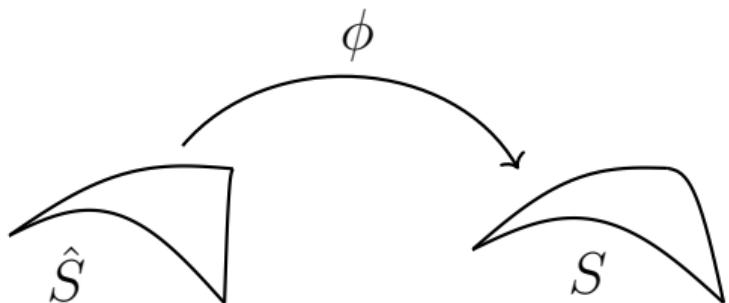
- $\mathbf{F} = \nabla_{\hat{\tau}} \phi$ ,  $J = \|\text{cof}(\mathbf{F})\|_F$



- Normal vector  $\hat{\nu}$
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- Element normal vector  $\hat{\mu} = \hat{\nu} \times \hat{\tau}$



- $\mathbf{F} = \nabla_{\hat{\tau}} \phi$ ,  $J = \|\text{cof}(\mathbf{F})\|_F$
- $\nu \circ \phi = \frac{1}{J} \text{cof}(\mathbf{F}) \hat{\nu}$
- $\tau \circ \phi = \frac{1}{J_B} \mathbf{F} \hat{\tau}$
- $\mu \circ \phi = \nu \circ \phi \times \tau \circ \phi$



$$\mathcal{W}(u) = \frac{t}{2} \|\boldsymbol{E}(u)\|_{\mathbb{M}}^2 + \frac{t^3}{24} \|\boldsymbol{F}^T \nabla(\nu \circ \phi) - \nabla \hat{\nu}\|_{\mathbb{M}}^2$$

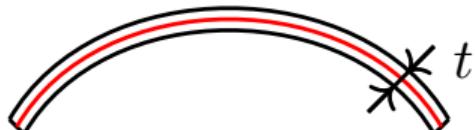
$u$  ... displacement of mid-surface

$t$  ... thickness

$\mathbb{M}$  ... material tensor

$$\boldsymbol{F} = \nabla u + \boldsymbol{P} = \nabla \phi, \quad \boldsymbol{P} = \boldsymbol{I} - \hat{\nu} \otimes \hat{\nu}$$

$$\boldsymbol{E} = \frac{1}{2}(\boldsymbol{F}^T \boldsymbol{F} - \boldsymbol{P}) = \frac{1}{2}(\nabla u^T \nabla u + \nabla u^T \boldsymbol{P} + \boldsymbol{P} \nabla u)$$



# Koiter shell

$$\mathcal{W}(u) = \frac{t}{2} \|\boldsymbol{E}(u)\|_{\mathbb{M}}^2 + \frac{t^3}{24} \|\boldsymbol{F}^T \nabla(\nu \circ \phi) - \nabla \hat{\nu}\|_{\mathbb{M}}^2$$

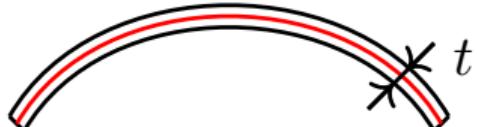
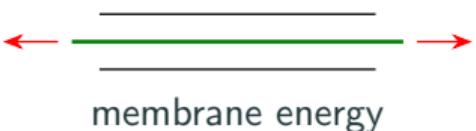
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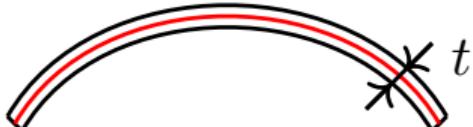
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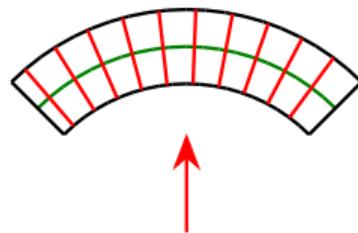
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membrane energy



bending energy

$$\mathcal{W}(u, \gamma) = \frac{t}{2} \|\boldsymbol{\mathcal{E}}(u)\|_{\mathbb{M}}^2 + \frac{t^3}{24} \|\text{sym}(\boldsymbol{F}^T \nabla(\tilde{\nu} \circ \phi)) - \nabla \hat{\nu}\|_{\mathbb{M}}^2 + \frac{t \kappa G}{2} \|\boldsymbol{F}^T \tilde{\nu} \circ \phi\|^2$$

$\gamma$  ... shearing

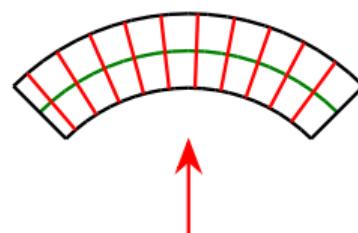
$$\tilde{\nu} = \frac{\nu + \gamma}{\|\nu + \gamma\|} \dots \text{director}$$

$G$  ... shearing modulus

$\kappa = 5/6$  ... shear correction factor



membrane energy



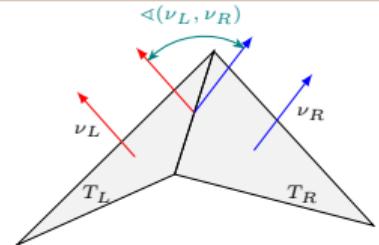
bending energy



shearing energy

# Distributional curvature for Koiter shell

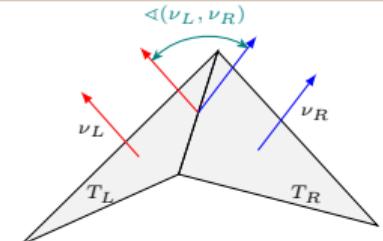
Lifting:  $\int_{\mathcal{T}_h} \kappa : \boldsymbol{\sigma} dx = \sum_{T \in \mathcal{T}_h} \int_T \nabla \nu : \boldsymbol{\sigma} dx + \sum_{E \in \mathcal{E}_h} \triangleleft(\nu_L, \nu_R) \boldsymbol{\sigma}_{\mu\mu} ds$



N., SCHÖBERL, STURM, Numerical shape optimization of Canham-Helfrich-Evans bending energy, *J. Comput. Phys.* (2023).

# Distributional curvature for Koiter shell

Lifting:  $\int_{\mathcal{T}_h} \kappa : \boldsymbol{\sigma} dx = \sum_{T \in \mathcal{T}_h} \int_T \nabla \nu : \boldsymbol{\sigma} dx + \sum_{E \in \mathcal{E}_h} \triangleleft(\nu_L, \nu_R) \boldsymbol{\sigma}_{\mu\mu} ds$



- Lifted curvature difference  $\kappa^{\text{diff}}$  via three-field Hu–Washizu formulation

$$\begin{aligned} \mathcal{L}(u, \kappa^{\text{diff}}, \boldsymbol{\sigma}) = & \frac{t}{2} \|\boldsymbol{E}(u)\|_{\mathbb{M}}^2 + \frac{t^3}{12} \|\kappa^{\text{diff}}\|_{\mathbb{M}}^2 - \langle f, u \rangle + \sum_{T \in \mathcal{T}_h} \int_T (\kappa^{\text{diff}} - (\boldsymbol{F}^T \nabla(\nu \circ \phi) - \nabla \hat{\nu})) : \boldsymbol{\sigma} dx \\ & + \sum_{E \in \mathcal{E}_h} \int_E (\triangleleft(\nu_L, \nu_R) - \triangleleft(\hat{\nu}_L, \hat{\nu}_R)) \boldsymbol{\sigma}_{\hat{\mu}\hat{\mu}} ds \end{aligned}$$

- Lagrange parameter  $\boldsymbol{\sigma} \in M_h^{k-1}$  moment tensor
- Eliminate  $\kappa^{\text{diff}} \rightarrow$  two-field formulation in  $(u, \boldsymbol{\sigma})$

N., SCHÖBERL: The Hellan–Herrmann–Johnson and TDNNS method for linear and nonlinear shells, arXiv:2304.13806.

## Shell problem

Find  $u \in [\mathcal{L}_h^k(\mathcal{T}_h)]^3$  and  $\sigma \in M_h^{k-1}$  for ( $H_\nu := \sum_i (\nabla^2 u_i) \nu_i$ )

$$\begin{aligned}\mathcal{L}(u, \sigma) = & \frac{t}{2} \|\boldsymbol{E}(u)\|_{\mathbb{M}}^2 - \frac{6}{t^3} \|\boldsymbol{\sigma}\|_{\mathbb{M}^{-1}}^2 - \langle f, u \rangle \\ & + \sum_{T \in \mathcal{T}_h} \int_T \boldsymbol{\sigma} : (H_\nu + (1 - \hat{\nu} \cdot \nu) \nabla \hat{\nu}) \, dx \\ & + \sum_{E \in \mathcal{E}_h} \int_E (\triangle(\nu_L, \nu_R) - \triangle(\hat{\nu}_L, \hat{\nu}_R)) \boldsymbol{\sigma}_{\hat{\mu} \hat{\mu}} \, ds\end{aligned}$$

Use hybridization to eliminate  $\sigma \rightarrow$  recover minimization problem

-  N., SCHÖBERL: The Hellan–Herrmann–Johnson method for nonlinear shells, *Comput. Struct.* 225 (2019).

$$H^1(\Omega) = \{u \in L^2(\Omega) \mid \nabla u \in [L^2(\Omega)]^d\}$$

$$\mathcal{L}_h^k(\mathcal{T}_h) = \mathcal{P}^k(\mathcal{T}_h) \cap C(\Omega)$$

$$H(\text{curl}, \Omega) = \{\boldsymbol{\sigma} \in [L^2(\Omega)]^d \mid \text{curl} \boldsymbol{\sigma} \in [L^2(\Omega)]^{2d-3}\}$$

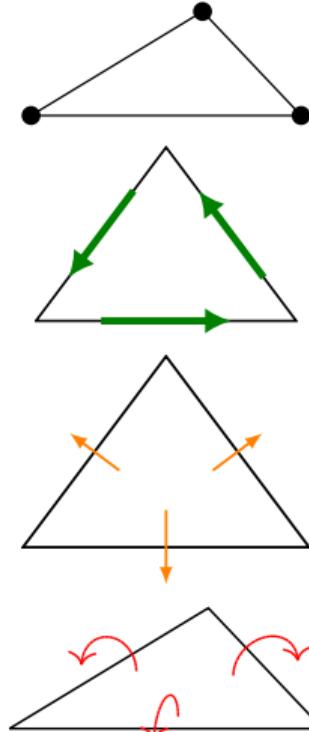
$$\mathcal{N}_{II}^k = \{\boldsymbol{\sigma} \in [\mathcal{P}^k(\mathcal{T}_h)]^d \mid [\![\boldsymbol{\sigma}_\tau]\!]_F = 0\}$$

$$H(\text{div}, \Omega) = \{\boldsymbol{\sigma} \in [L^2(\Omega)]^d \mid \text{div} \boldsymbol{\sigma} \in L^2(\Omega)\}$$

$$BDM^k = \{\boldsymbol{\sigma} \in [\mathcal{P}^k(\mathcal{T}_h)]^d \mid [\![\boldsymbol{\sigma}_n]\!]_F = 0\}$$

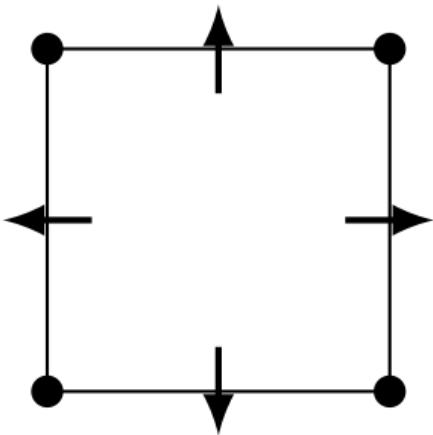
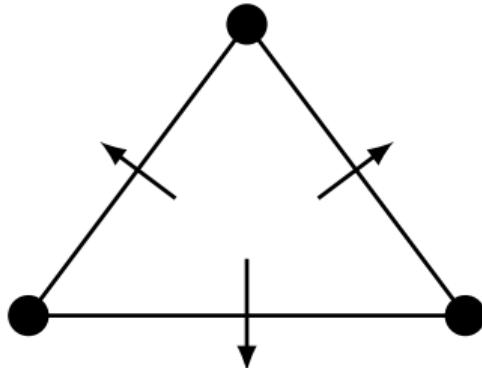
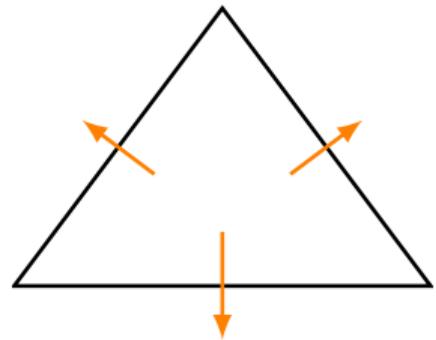
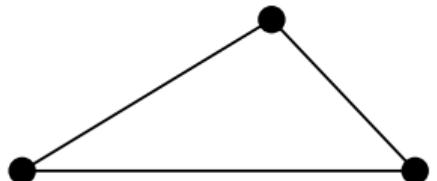
$$H(\text{divdiv}, \Omega) = \{\boldsymbol{\sigma} \in [L^2(\Omega)]_{\text{sym}}^{d \times d} \mid \text{divdiv} \boldsymbol{\sigma} \in H^{-1}(\Omega)\}$$

$$M_h^k(\mathcal{T}_h) = \{\boldsymbol{\sigma} \in [\mathcal{P}^k(\mathcal{T}_h)]_{\text{sym}}^{d \times d} \mid [\![\mathbf{n}^T \boldsymbol{\sigma} \mathbf{n}]\!]_F = 0\}$$



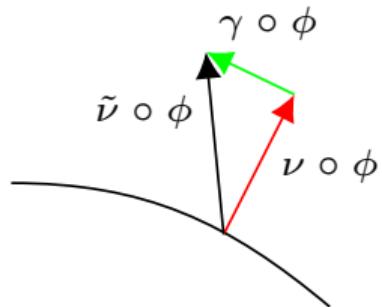
- A. PECHSTEIN AND J. SCHÖBERL: The TDNNS method for Reissner-Mindlin plates, *J. Numer. Math.* (2017) 137, pp. 713-740.

# Shell element (Koiter)



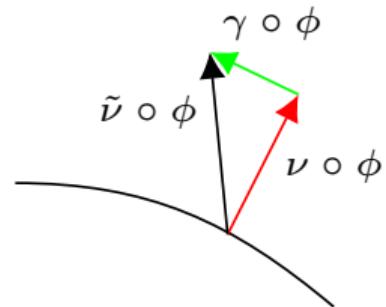
# Extension to nonlinear Naghdi shells

- Use hierarchical shell model
- Additional shearing dofs  $\gamma$  in  $H(\text{curl})$
- $\tilde{\nu} \circ \phi = \frac{\nu \circ \phi + \gamma \circ \phi}{\|\nu \circ \phi + \gamma \circ \phi\|}$
- Free of shear locking



 ECHTER, R. AND OESTERLE, B. AND BISCHOFF, M.: A hierachic family of isogeometric shell finite elements, *Comput. Methods Appl. Mech. Engrg* (2013) 254, pp. 170–180.

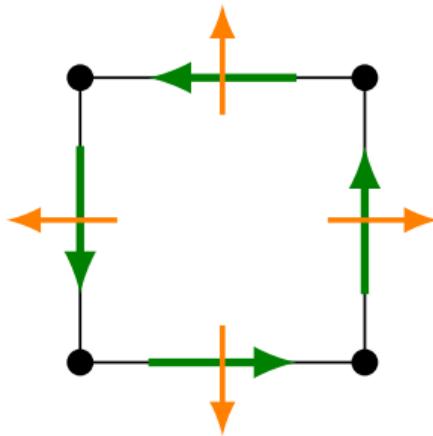
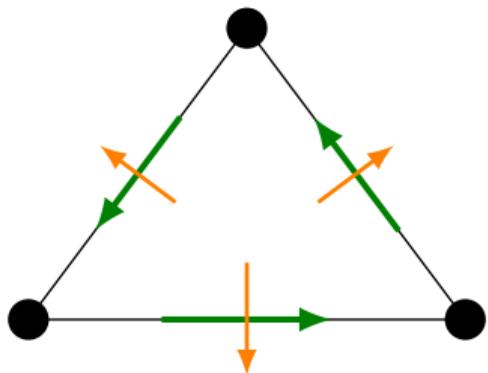
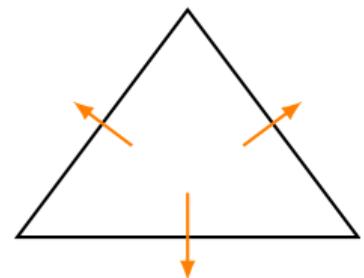
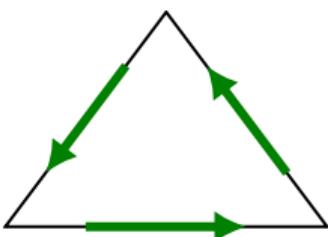
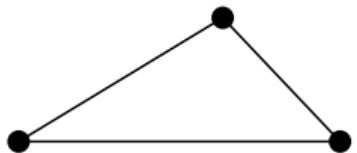
- Use hierarchical shell model
- Additional shearing dofs  $\gamma$  in  $H(\text{curl})$
- $\tilde{\nu} \circ \phi = \nu \circ \phi + \gamma \circ \phi = \frac{1}{J} \text{cof}(\boldsymbol{F}) \hat{\nu} + (\boldsymbol{F}^\dagger)^\top \hat{\gamma}$
- Free of shear locking



$$\begin{aligned}\mathcal{L}(u, \sigma, \hat{\gamma}) &= \frac{t}{2} \|\boldsymbol{E}(u)\|_{\mathbb{M}}^2 + \frac{t\kappa G}{2} \|\hat{\gamma}\|^2 - \frac{6}{t^3} \|\sigma\|_{\mathbb{M}^{-1}}^2 \\ &\quad + \sum_{T \in \mathcal{T}_h} \int_T (\boldsymbol{H}_{\tilde{\nu}} + (1 - \tilde{\nu} \cdot \hat{\nu}) \nabla \hat{\nu} - \nabla \hat{\gamma}) : \sigma \, dx \\ &\quad + \sum_{E \in \mathcal{E}_h} \int_E (\llangle (\nu_L, \nu_R) - \llangle (\hat{\nu}_L, \hat{\nu}_R) + [\![ \hat{\gamma}_{\hat{\mu}} ]\!] ) \sigma_{\hat{\mu} \hat{\mu}} \, ds\end{aligned}$$

 N., SCHÖBERL: The Hellan–Herrmann–Johnson and TDNNS method for linear and nonlinear shells, [arXiv:2304.13806](https://arxiv.org/abs/2304.13806).

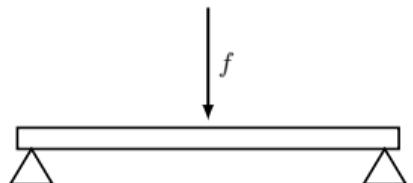
# Shell element (Naghdi)



$$\mathcal{L}_{\text{lin}}^{\text{shell}}(u, \boldsymbol{\sigma}) = \frac{t}{2} \|\text{sym}(\nabla^{\text{cov}} u)\|_{\mathbb{M}}^2 - \frac{6}{t^3} \|\boldsymbol{\sigma}\|_{\mathbb{M}^{-1}}^2 + \sum_{T \in \mathcal{T}_h} \left( \int_T \boldsymbol{H}_{\hat{\nu}} : \boldsymbol{\sigma} dx - \int_{\partial T} (\nabla u^\top \hat{\nu})_{\hat{\mu}} \boldsymbol{\sigma}_{\hat{\mu}\hat{\mu}} ds \right)$$

$$\mathcal{L}_{\text{lin}}^{\text{plate}}(w, \boldsymbol{\sigma}) = -\frac{6}{t^3} \|\boldsymbol{\sigma}\|_{\mathbb{M}^{-1}}^2 + \sum_{T \in \mathcal{T}_h} \left( \int_T \nabla^2 w : \boldsymbol{\sigma} dx - \int_{\partial T} \frac{\partial w}{\partial \hat{\mu}} \boldsymbol{\sigma}_{\hat{\mu}\hat{\mu}} ds \right)$$

$$\text{divdiv} \nabla^2 w = f \Leftrightarrow \begin{cases} \boldsymbol{\sigma} = \nabla^2 w, \\ \text{divdiv} \boldsymbol{\sigma} = f, \end{cases}$$



 M. COMODI: The Hellan-Herrmann-Johnson method: some new error estimates and postprocessing, *Math. Comp.* 52 (1989) pp. 17–29.

$$\begin{aligned}\mathcal{L}_{\text{lin}}^{\text{shell}}(u, \boldsymbol{\sigma}, \hat{\boldsymbol{\gamma}}) &= \frac{t}{2} \|\text{sym}(\nabla^{\text{cov}} u)\|_{\mathbb{M}}^2 + \frac{t\kappa G}{2} \|\hat{\boldsymbol{\gamma}}\|^2 - \frac{6}{t^3} \|\boldsymbol{\sigma}\|_{\mathbb{M}^{-1}}^2 \\ &\quad + \sum_{T \in \mathcal{T}_h} \left( \int_T (\boldsymbol{H}_{\hat{\nu}} - \nabla \hat{\boldsymbol{\gamma}}) : \boldsymbol{\sigma} \, dx - \int_{\partial T} ((\nabla u^\top \hat{\nu})_{\hat{\mu}} - \hat{\boldsymbol{\gamma}}_{\hat{\mu}}) \boldsymbol{\sigma}_{\hat{\mu}\hat{\mu}} \, ds \right) \\ \mathcal{L}_{\text{lin}}^{\text{plate}}(w, \boldsymbol{\sigma}, \hat{\boldsymbol{\gamma}}) &= \frac{t\kappa G}{2} \|\hat{\boldsymbol{\gamma}}\|^2 - \frac{6}{t^3} \|\boldsymbol{\sigma}\|_{\mathbb{M}^{-1}}^2 \\ &\quad + \sum_{T \in \mathcal{T}_h} \left( \int_T (\nabla^2 w - \nabla \hat{\boldsymbol{\gamma}}) : \boldsymbol{\sigma} \, dx - \int_{\partial T} \left( \frac{\partial w}{\partial \hat{\mu}} - \hat{\boldsymbol{\gamma}}_{\hat{\mu}} \right) \boldsymbol{\sigma}_{\hat{\mu}\hat{\mu}} \, ds \right)\end{aligned}$$

-  A. PECHSTEIN AND J. SCHÖBERL: The TDNNS method for Reissner–Mindlin plates, *J. Numer. Math.* (2017) 137, pp. 713–740.

## Membrane locking

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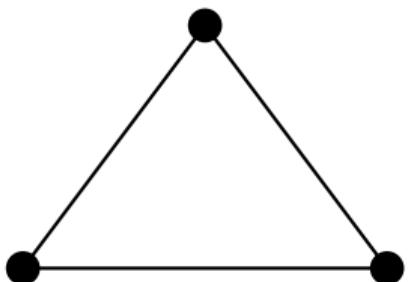
$$\mathcal{W}(u) = t E_{\text{mem}}(u) + t^3 E_{\text{bend}}(u) - f \cdot u, \quad f = t^3 \tilde{f}$$

$$\mathcal{W}(u) = t^{-2} E_{\text{mem}}(u) + E_{\text{bend}}(u) - \tilde{f} \cdot u, \quad f = t^3 \tilde{f}$$

Enforces  $E_{\text{mem}}(u) = 0$  in the limit  $t \rightarrow 0$

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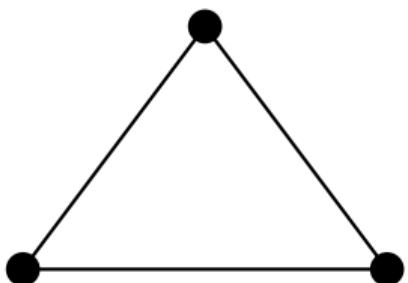


$$\mathcal{L}_h^k(\mathcal{T}_h) = \mathcal{P}^k(\mathcal{T}_h) \cap C(\Omega) \subset H^1(\Omega)$$

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$$E_{\text{mem}}(u) = 0 \quad \not\Rightarrow \quad E_{\text{mem}}(u_h) = 0$$

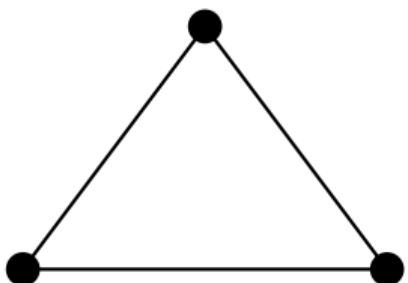


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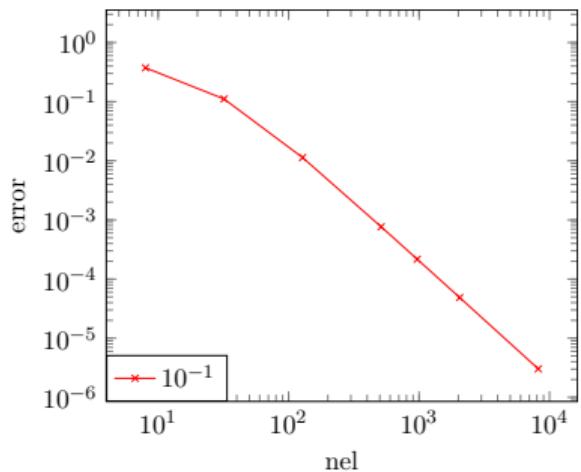
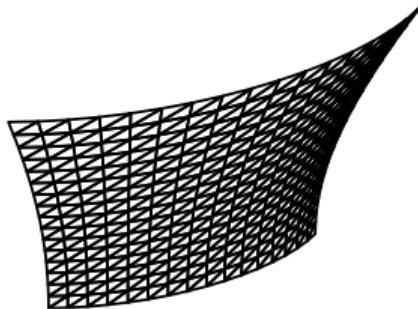
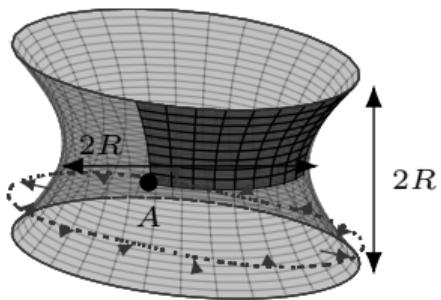
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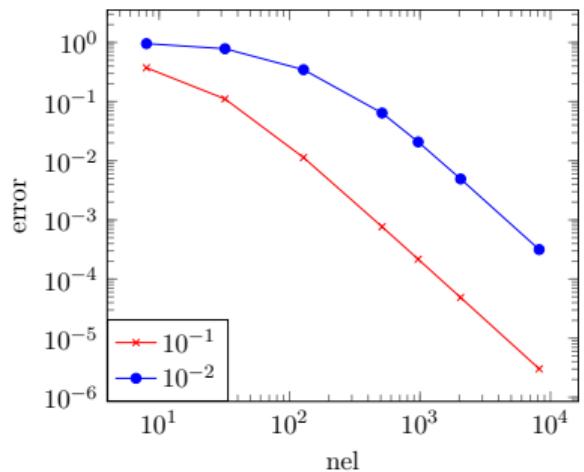
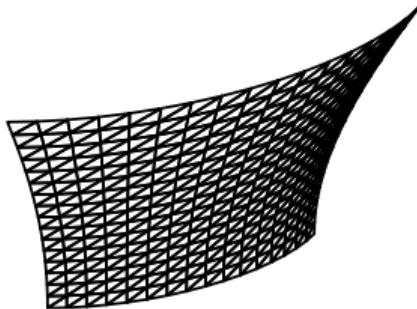
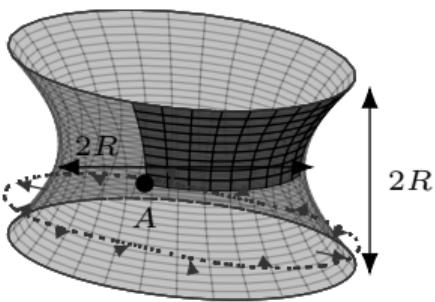


$$\mathcal{L}_h^k(\mathcal{T}_h) = \mathcal{P}^k(\mathcal{T}_h) \cap C(\Omega) \subset H^1(\Omega)$$

# Hyperboloid with free ends

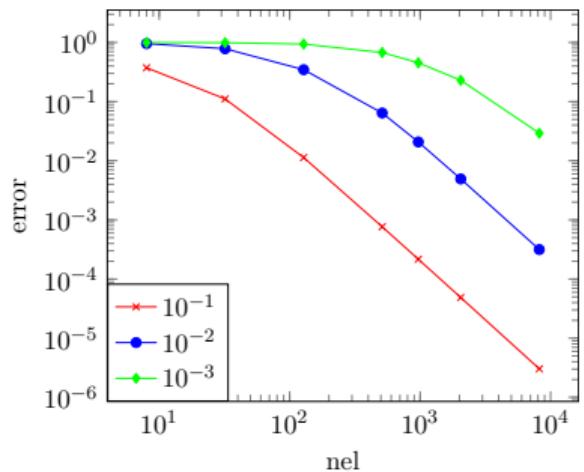
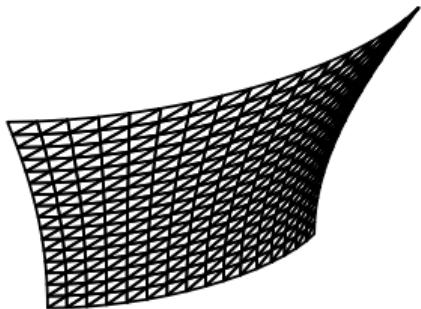
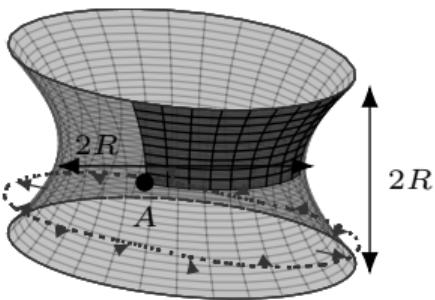


# Hyperboloid with free ends



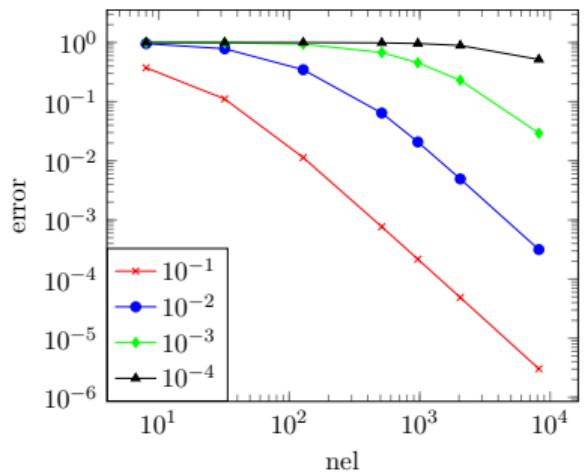
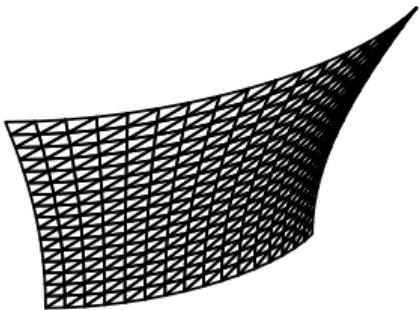
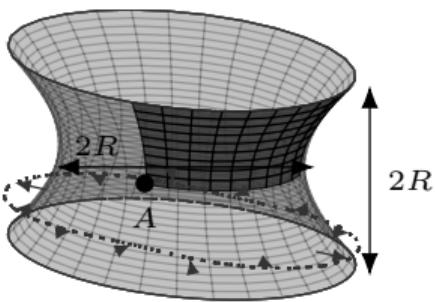
- Pre-asymptotic regime

# Hyperboloid with free ends



- Pre-asymptotic regime

# Hyperboloid with free ends

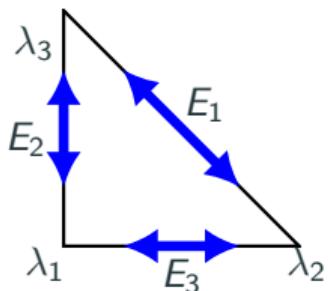


- Pre-asymptotic regime

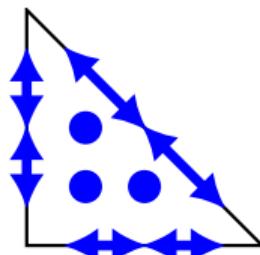
# Regge elements

$$H(\operatorname{curl} \operatorname{curl}) := \{\boldsymbol{\sigma} \in [L^2(\Omega)]_{\text{sym}}^{2 \times 2} \mid \operatorname{curl} \operatorname{curl} \boldsymbol{\sigma} \in H^{-1}(\Omega)\}$$

$$\operatorname{Reg}_h^k := \{\boldsymbol{\varepsilon} \in [\mathcal{P}^k(\mathcal{T}_h)]_{\text{sym}}^{d \times d} \mid [\![t^\top \boldsymbol{\varepsilon} t]\!]_E = 0 \text{ for all edges } E\}$$



$$\varphi_{E_i} = \nabla \lambda_j \odot \nabla \lambda_k, \quad t_j^\top \varphi_{E_i} t_j = c_i \delta_{ij},$$

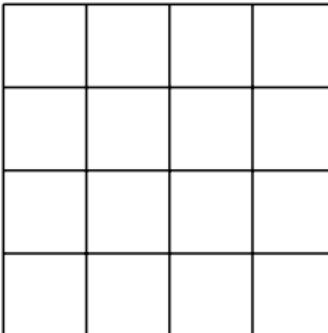


$$\varphi_{T_i} = \lambda_i \nabla \lambda_j \odot \nabla \lambda_k$$

-  CHRISTIANSEN: On the linearization of Regge calculus, *Numerische Mathematik* 119, 4 (2011).
-  LI: Regge Finite Elements with Applications in Solid Mechanics and Relativity, *PhD thesis, University of Minnesota* (2018).
-  N.: Mixed Finite Element Methods For Nonlinear Continuum Mechanics And Shells, *PhD thesis, TU Wien* (2021).

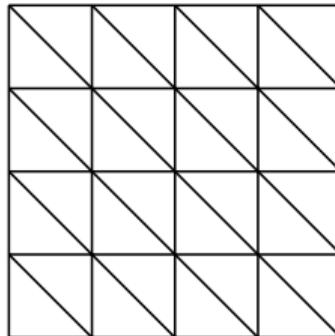
$$\frac{1}{t^2} \| \mathbf{E}(u_h) \|_{\mathbb{M}}^2$$

$$\frac{1}{t^2} \|\Pi_{L^2}^k E(u_h)\|_{\mathbb{M}}^2$$

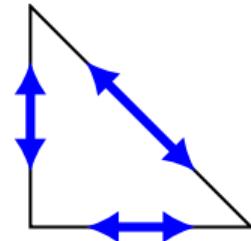


- Reduced integration for quadrilateral meshes

$$\frac{1}{t^2} \|\mathcal{I}_{\mathcal{R}}^k E(u_h)\|_{\mathbb{M}}^2$$

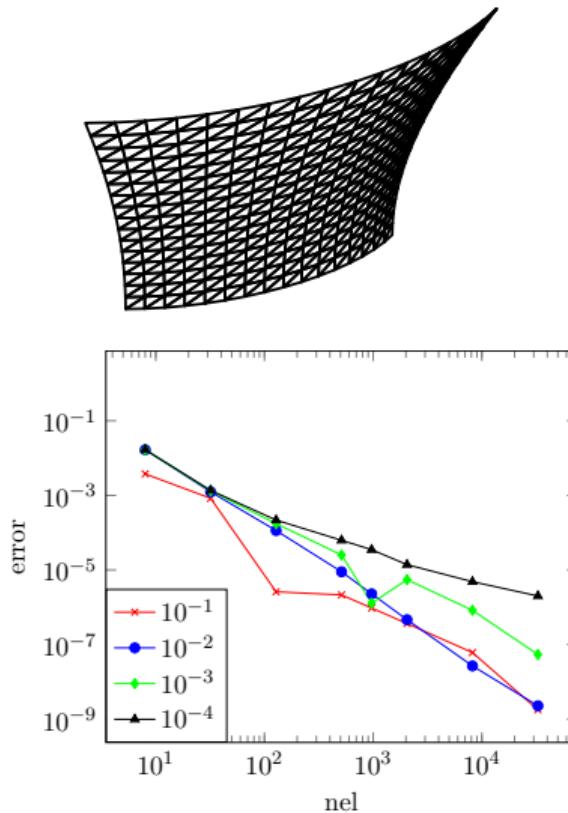
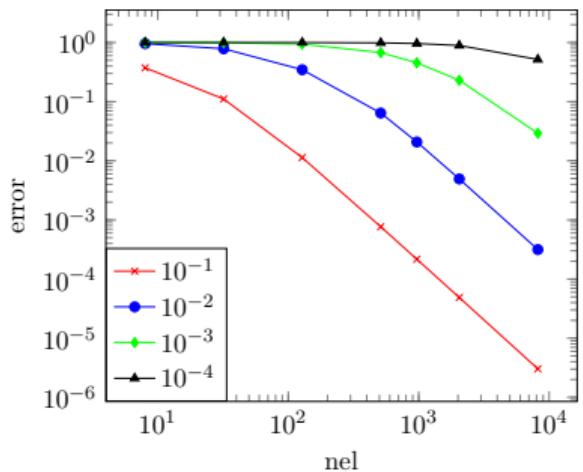
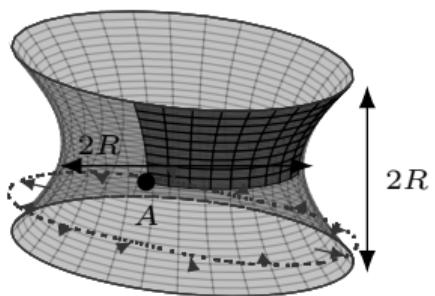


- Reduced integration for quadrilateral meshes
- Regge interpolant for triangles
- Connection to MITC shell elements



 N., SCHÖBERL: Avoiding membrane locking with Regge interpolation, *Comput. Methods Appl. Mech. Engrg* 373 (2021).

# Hyperboloid with free ends



## Numerical examples

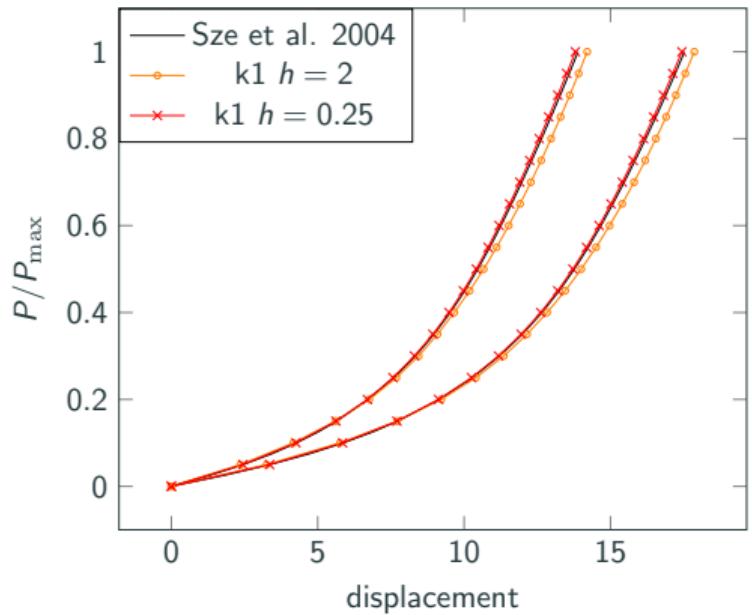
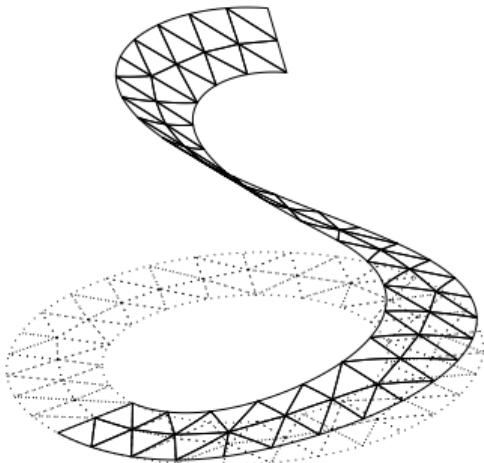
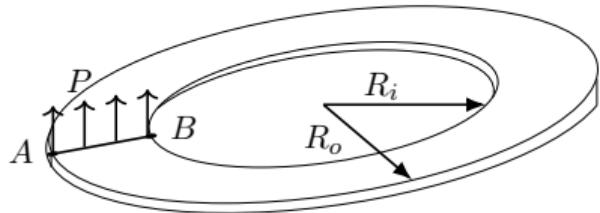
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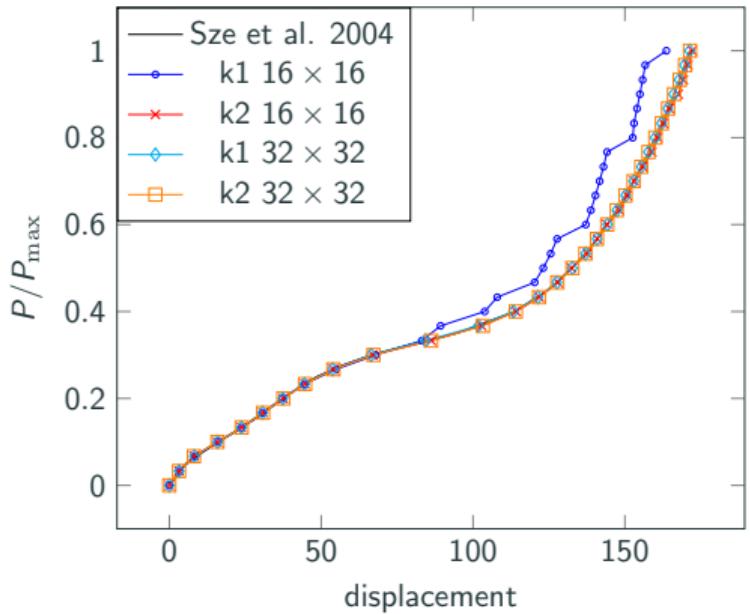
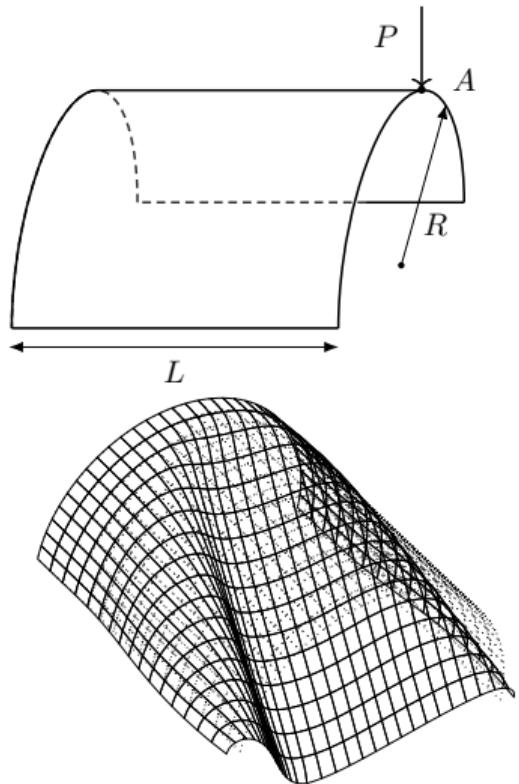
# Cantilever subjected to end moment

# Cantilever subjected to end moment

# Slit annular plate



# Pinched cylinder



- Lifting of distributional curvature via Hu–Washizu three-field formulation
- HHJ for (non)linear Koiter shells
- TDNNS for (non)linear Naghdi shells
- Avoiding membrane locking with Regge elements

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- 
- Coupling for 3D elasticity (A. Pechstein, M. Krommer; JKU)
  - NGSolve Add-On

-  N., SCHÖBERL: The Hellan–Herrmann–Johnson and TDNNS method for linear and nonlinear shells, *arXiv:2304.13806*.
-  N., SCHÖBERL: The Hellan–Herrmann–Johnson method for nonlinear shells, *Comput. Struct.* 225 (2019).
-  N., SCHÖBERL: Avoiding membrane locking with Regge interpolation, *Comput. Methods Appl. Mech. Engrg* 373 (2021).
-  N.: Mixed Finite Element Methods For Nonlinear Continuum Mechanics And Shells, *PhD thesis, TU Wien* (2021).
-  N., SCHÖBERL, STURM, Numerical shape optimization of Canham-Helfrich-Evans bending energy, *J. Comput. Phys.* (2023).

**Thank You for Your attention!**